



## Influence of Wavelengths on the Photocurrent Density and on the Internal Quantum Efficiency of an Irradiated Monofacial Photopile of Type N + PP +

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### ABSTRACT

In this article, the influence of the wavelengths of electromagnetic radiation on the photocurrent density and on the percentage of the number of electrons created and collected relative to the number of photons absorbed from an irradiated monofacial solar cell has made the object of the study.

From the establishment of the scattering coefficient and the density of the minority carriers, the expressions of the photocurrent density and the internal quantum efficiency are established and studied as a function of the wavelength for different: angular frequencies, energies of irradiations and damage coefficients.

**Key words:** Irradiating particles, Photocell, wavelength, Internal Quantum Efficiency, UV, Infrared

### INTRODUCTION

The solar spectrum used as a reference in the photovoltaic industry corresponds to an illumination of at a temperature of with a crossing (Air Mass where G means global because taking into account both direct and diffuse radiation) called the standard test conditions (STC) [1]. This definition of the spectrum is all the more important as it will impact the performance of photovoltaic devices. The limited performance of solar cells has been studied by several groups of researchers including the famous work of Shockley and Queisser [2]. There is also the work of M. Green [3], H. Kiess [4] and F. Meillaud et al [5]. The various loss factors are due to purely physical restrictions related to the material, or to technological limitations induced by the manufacturing process.

1) Physical losses

- i. Losses by long wavelength photons
- ii. Losses due to excess photon energy
- iii. Voltage factor
- iv. Fillfactor

2) Technological losses

- I. Reflectivity
- ii. Shading rate
- iii. Absorption efficiency
- iv. Collection efficiency

In addition to these limitations, there is also the influence of the series resistance and the shunt resistance.

The fill factor ( $ff$ ) and the efficiency ( $\eta$ ) provide information on the overall performance of the structure. To detail these characteristics as a function of the wavelength used, the parameters  $IQE$ ;  $EQE$  (respectively internal quantum efficiency and external quantum efficiency) are introduced.

In this article, the variation of the photocurrent density and the internal quantum efficiency ( $IQE$ ) of an irradiated monofacial solar cell, under monochromatic illumination and under frequency condition is highlighted.

STUDYMODEL

Study device

The following figure represents the single-face silicon-based study solar cell of the type  $n^+pp^+$  [6],[7], initially irradiated with charged particles and subjected to monochromatic illumination.

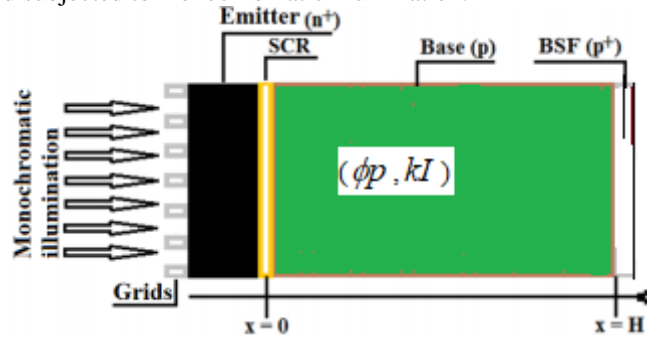


Fig. 1 An  $n^+pp^+$  type of a silicon solar cell under irradiation

The continuity equation that governs the phenomena that take place there after illumination is:

$$D(\omega, kI, \phi p, T) \frac{\partial^2 \delta(x, t)}{\partial x^2} - \frac{\delta(x, t)}{\tau} + g(x, t) = + \frac{\partial \delta(x, t)}{\partial t} \tag{1}$$

Where,

i-  $D(\omega, kI, \phi p, T)$  represents the complex diffusion coefficient of electrons in the base of the solar cell under irradiation and under temperature [8], [9], [10],[11]. F. Ba et al establish its full expression with an explicit determination of the parameters in the article [12] where the temperature 310 K corresponds to a diffusion coefficient of which is specific to microcrystalline silicon.

ii-  $\delta(x, t)$  and  $g(x, t)$  respectively represent the density of the minority charge carriers and the rate of generation in white light of the excess charge carriers as a function of the thickness of the base and of the modulation frequency [13], [14]:

$$\delta(x, t) = \delta(x) \cdot \exp(j\omega t) \tag{2}$$

$$g(x, t) = g(x) \cdot \exp(j\omega t) \tag{3}$$

$$g(x) = \alpha_\lambda \cdot \phi_\lambda \cdot (1 - R_\lambda) \cdot \exp(-\alpha_\lambda \cdot x) \tag{4}$$

$\alpha_\lambda$  represents the monochromatic absorption coefficient of the material,  $R_\lambda$ : the monochromatic reflection coefficient of the material for a wavelength  $\lambda$ ;  $\phi_\lambda$ : the incident flux of monochromatic light; the thickness of the base of the solar cell or the depth of penetration of the incident flux of light.

Solving the equation gives the expression for the density of minority carriers written as follows:

$$\delta(x, \omega, \lambda, kI, \phi p, T) = A \cdot \cosh\left(\frac{x}{L(\omega, kI, \phi p, T)}\right) + B \cdot \sinh\left(\frac{x}{L(\omega, kI, \phi p, T)}\right) + K(\omega, kI, \phi p, T) \cdot \exp(-\alpha_\lambda \cdot x) \tag{5}$$

With,

$$K(\omega, \lambda, kI, \phi p, T) = - \frac{\alpha_\lambda \cdot \phi_\lambda \cdot (1 - R_\lambda) \cdot L(\omega, kI, \phi p, T)^2}{D(\omega, kI, \phi p, T) \cdot [\alpha_\lambda^2 L_\omega^2 - 1]} \tag{6}$$

The coefficients A and B will be determined from the boundary conditions [15]:

➤ At the base transmitter junction

$$\frac{\partial \delta(x, \omega, \lambda, kI, \phi p, T)}{\partial x} \Big|_{x=0} = \frac{Sf}{D(\omega, kI, \phi p, T)} \delta(0, \omega, \lambda, kI, \phi p, T) \tag{7}$$

$Sf$ : represents the rate of recombination at the junction [16],[17],[18]

$$Sf = Sf_o + Sf_m \tag{8}$$

$Sf_o$  : rate of recombination intrinsically induced by the shunt resistance and depending only on the intrinsic parameters of the solar cell

$Sf_m = m \cdot 10^m$  : Recombination rate which relates to the external load imposing the operating point of the solar cell

➤ On the back side

$$\left. \frac{\partial \delta(x, \omega, \lambda, kI, \phi p, T)}{\partial x} \right|_{x=H} = - \frac{Sb}{D(\omega, kI, \phi p, T)} \delta(H, \omega, \lambda, kI, \phi p, T) \tag{9}$$

$Sb$  : Recombination rate of charge carriers on the rear face [19], [20].

The internal quantum efficiency ( $IQE$ ) is the number of electrons created and collected for each absorbed photon. It is obtained by taking into account the specific characteristics of the cell (diffusion length, surface and volume recombination); its expression is given as follows:

$$IQE(\omega, j, p, \lambda, kI, \phi p, T) = \frac{EQE(\omega, j, p, \lambda, kI, \phi p, T)}{(1 - R_\lambda)} \tag{10}$$

$(1 - R_\lambda)$  : is the transmitted part of the incident photon flux in the material.

$$EQE(\omega, j, p, \lambda, kI, \phi p, T) = \frac{Jph(\omega, j, p, \lambda, kI, \phi p, T)}{q \cdot \phi_\lambda} \tag{11}$$

$$IQE(\omega, j, p, \lambda, kI, \phi p, T) = \frac{Jph(\omega, j, p, \lambda, kI, \phi p, T)}{q \cdot (1 - R_\lambda) \cdot \phi_\lambda} \tag{12}$$

$q$  : Elementary charge of the electron;  $Jph(\omega, j, p, \lambda, kI, \phi p, T)$  : photocurrent density, its expression is given by the relation called Fick relation.

$$Jph(\omega, j, p, \lambda, kI, \phi p, T) = q \cdot D(\omega, kI, \phi p, T) \cdot \left. \frac{\partial \delta(x, \omega, \lambda, kI, \phi p, T)}{\partial x} \right|_{x=0} \tag{13}$$

The figure below illustrates the variation of the photocurrent density as a function of the wavelength for different values of the rate of recombination at the junction.

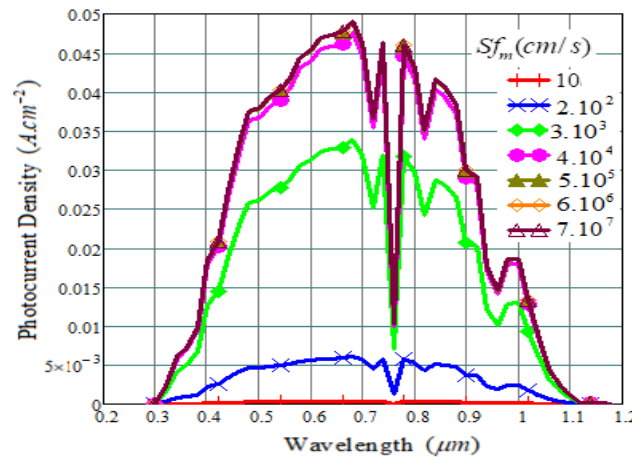


Fig. 2 Photocurrent Density versus Wavelength for different values of the recombination velocity at the junction.

$$\omega = 10^5 \text{ rad/s} ; kI = 15 \text{ cm}^{-2} / \text{MeV} ; \phi p = 50 \text{ MeV} ; T = 310 \text{ K}$$

From this plot, we have a small variation in the photocurrent density for low and long wavelengths. The absorption peak is obtained at  $\lambda = 0.68 \mu\text{m}$  where the photocurrent density is at its maximum. Beyond that, there is a decrease in the latter as a function of the wavelength.

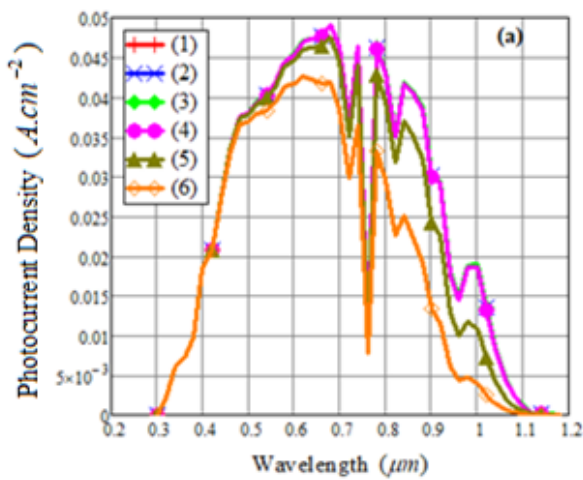
Only energy photons  $E = \frac{h \cdot c}{\lambda} > E_g = \frac{h \cdot c}{\lambda_g}$  ( $h = 6,62 \times 10^{-34} J.s$ : Planck constant;  $c = 3 \times 10^8 m.s^{-1}$ : speed of

light ;  $E_g$  : silicon gap energy) can be called. The absorbed energy allows electrons to pass from the valence band to the conduction band, hence the positive gradient observed for the photocurrent density at low wavelengths (part of the visible + UV).

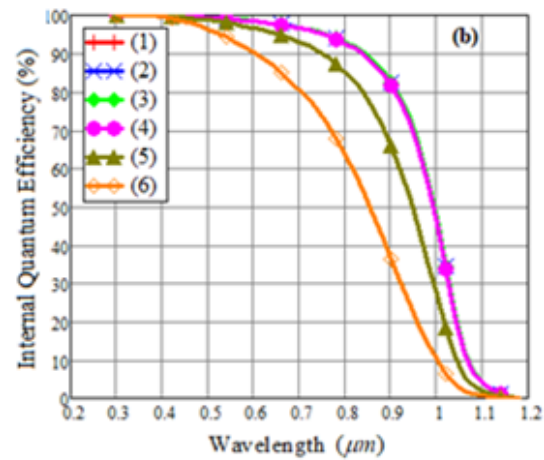
In the infrared range, increasing the wavelength decreases the photocurrent density. And this decrease is more marked for the high recombination speeds at the junction. The following table shows a distribution of the wavelengths and the gap energy of silicon at a temperature  $T = 300K$

**Table-1** Wavelength of electromagnetic radiation

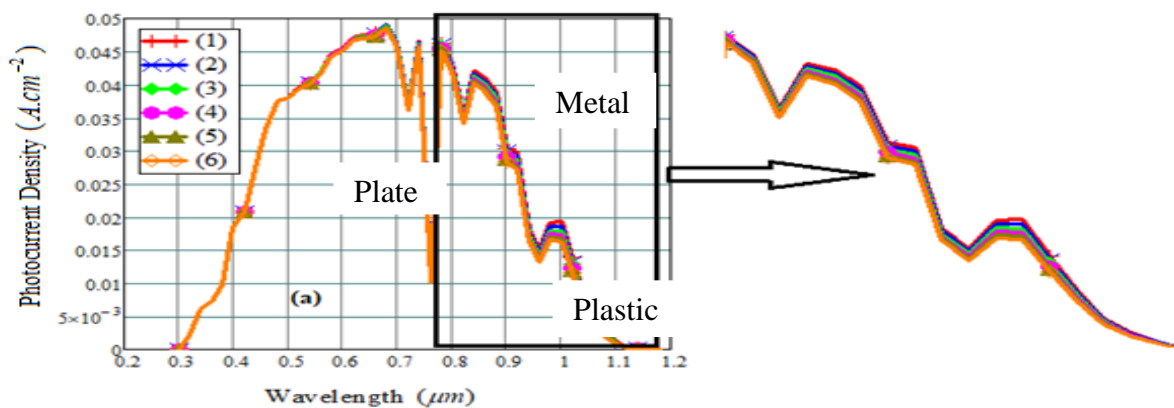
Ray	Wavelength	Photon Energy
UV	100 – 380 nm	3.26 – 124 eV
VISIBLE	380 – 750 nm	1.65 – 3.26 eV
INFRARED	750 – 1200 nm	0.0012 – 1.65 eV



**Figure 3a**



**Figure 3b**



**Figure 4a**

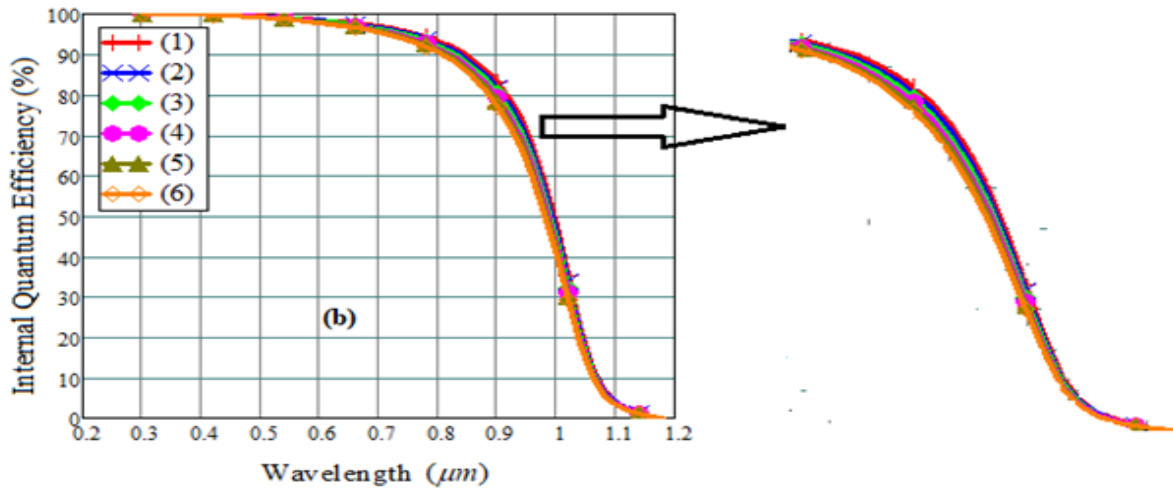


Figure 4b

Fig. 3,4,5: (a) Photocurrent Density and (b) Internal Quantum Efficiency versus Wavelength

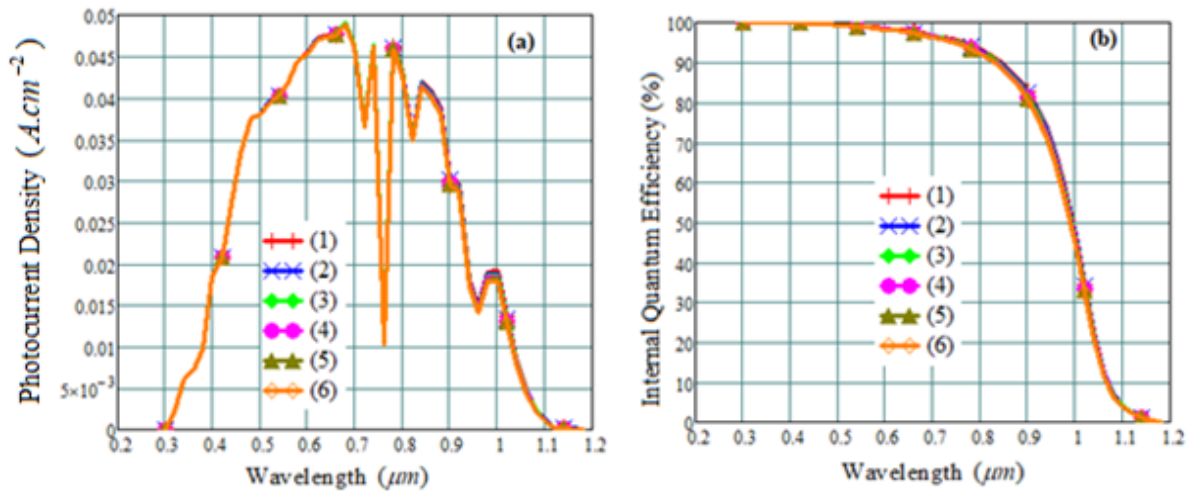


Figure 5a

Figure 5b

Fig. 3,4,5: (a) Photocurrent Density and (b) Internal Quantum Efficiency versus Wavelength

Table -2 List of different data taken to plot the figures above:

Figures 3(a,b)	For $\neq \omega(\text{rad} / \text{s})$ (1): $10^2 \text{ rad} / \text{s}$ ; (2): $10^3 \text{ rad} / \text{s}$ ; (3): $10^4 \text{ rad} / \text{s}$ ; (4): $10^5 \text{ rad} / \text{s}$ ; (5): $10^6 \text{ rad} / \text{s}$ ; (6): $10^7 \text{ rad} / \text{s}$ ; $\phi p = 50 \text{ MeV}$ ; $kI = 15 \text{ cm}^{-2} / \text{ MeV}$ ;
Figures 4(a,b)	For $\neq \phi p(\text{MeV})$ (1): $0 \text{ MeV}$ ; (2): $50 \text{ MeV}$ ; (3): $100 \text{ MeV}$ ; (4): $150 \text{ MeV}$ ; (5): $200 \text{ MeV}$ ; (6): $250 \text{ MeV}$ ; $\omega = 10^5 \text{ rad} / \text{s}$
Figures 5(a,b)	For $\neq kI(\text{cm}^{-2} / \text{ MeV})$ (1): $5 \text{ cm}^{-2} / \text{ MeV}$ ; (2): $10 \text{ cm}^{-2} / \text{ MeV}$ ; (3): $15 \text{ cm}^{-2} / \text{ MeV}$ ; (4): $20 \text{ cm}^{-2} / \text{ MeV}$ ; (5): $25 \text{ cm}^{-2} / \text{ MeV}$ ; $Sf_m = 6.10^6 \text{ cm} / \text{s}$ ; $\omega = 10^5 \text{ rad} / \text{s}$
$T = 310 \text{ K}$ ; $Sf_m = 6.10^6 \text{ cm} / \text{s}$ ; $Sb(\omega, \lambda, kI, \phi p, T)$	

The various quantum efficiency profiles as a function of the wavelength have generally obtained a constancy and then a progressive decrease as a function of the wavelength.

- ❖ For  $0.3\mu\text{m} \leq \lambda \leq 0.68\mu\text{m}$ , the internal quantum efficiency leaves an absolute maximum and then decreases as the wavelength increases (from 99% to 97%). There is a strong absorption of very energetic incident photons.
- ❖ For  $0.68\mu\text{m} < \lambda \leq 1.12\mu\text{m}$ , the internal quantum efficiency gradually decreases (from 97% to 2%). At these wavelengths, the absorption coefficient of silicon becomes slower and lower, resulting in low absorption of less energetic incident photons.

Under the effect of the angular frequency and of the irradiating particles (irradiation energy and damage coefficient  $kI$ ), the amplitude of the photocurrent density and of the internal quantum efficiency decreases. We have :

- under the effect of the angular frequency ( $\omega$ ) : a break in the generation of excess minority carriers thus leading to a decrease in the photo-created carriers participating in the photocurrent.
- under the effect of the irradiation energy ( $\phi p$ ) and the damage coefficient ( $kI$ ) : a decrease in the diffusion coefficient of the minority carriers and consequently in the photocurrent density.

The diminutive because the energy of photons in the infrared domain is much lower than the energy of the band gap or gap energy. In addition, variations in internal quantum efficiency are related to the effects of back surface recombination and the short diffusion length of carriers. As the wavelength increases, these effects increase and the internal quantum efficiency decreases.

### CONCLUSION

In this article, the effect of wavelength on photocurrent density and on internal quantum efficiency has been investigated. Short wavelengths (part of Ultra Violet + Visible) are designated on the front panel, while long wavelengths (Infrared) can be designated on the rear panel. The different variations of the internal quantum efficiency are explained by the variation of the photocurrent. In addition, the internal quantum is a function of the photocurrent density.

Thus, the study of internal quantum efficiency provides information on the quality of the front and rear of our solar cell.

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